

# 零温与有限温格林函数的关系

$\Delta$  零温:  $G^{R/A} = \int_{-\infty}^{+\infty} \frac{dw'}{2\pi} \frac{A(w')}{w-w \pm i0_{\pm}}$   
 有限温:  $G(w_n) = \int_{-\infty}^{+\infty} \frac{dw'}{2\pi} \frac{A(w')}{i w_n - w'}$

$i w_n \rightarrow \begin{cases} w + i0_{+} (G^R) \\ w - i0_{+} (G^A) \end{cases}$  解析延拓

## 费曼图, 费曼规则, 微扰论

例: 电声相互作用

$$\hat{H}_I = g \int \underbrace{\psi_{\alpha}^{\dagger}(\vec{r}, \tau) \psi_{\alpha}(\vec{r}, \tau)}_{\text{电子密度算符}} \underbrace{\phi(\vec{r}, \tau)}_{\text{声子场算符}} d\vec{r}$$

声子场算符 (忽略了极化)

$$\phi(\vec{r}) = \frac{1}{\sqrt{V}} \sum_{\vec{k}} \sqrt{\frac{\omega(\vec{k})}{2}} (\hat{b}_{\vec{k}} e^{i\vec{k}\cdot\vec{r}} + \hat{b}_{\vec{k}}^{\dagger} e^{-i\vec{k}\cdot\vec{r}})$$

变到相互作用表象

$$\phi(\vec{r}, \tau) = \frac{1}{\sqrt{V}} \sum_{\vec{k}} \sqrt{\frac{\omega(\vec{k})}{2}} [\hat{b}_{\vec{k}} e^{i\vec{k}\cdot\vec{r} - i\omega(\vec{k})\tau} + \hat{b}_{\vec{k}}^{\dagger} e^{-i\vec{k}\cdot\vec{r} + i\omega(\vec{k})\tau}]$$

目的, 电子格林函数

$$G(\tau_1 > \tau_2) = -\frac{1}{\Xi} \text{tr} (e^{-\beta(\hat{H} - \mu\hat{N})} \underbrace{\hat{\psi}(\vec{r}_1, \tau_1) \hat{\psi}^{\dagger}(\vec{r}_2, \tau_2)}_{\hat{H} = \hat{H}_0 + \hat{H}_{int}})$$

$$\Rightarrow e^{\tau_1(\hat{H} - \mu\hat{N})} \hat{\psi}(\vec{r}_1) e^{-\tau_1(\hat{H} - \mu\hat{N})}$$

海森堡表象

复习 S 矩阵  $S(\tau)$

$$e^{-\tau(\hat{H} - \mu\hat{N})} = e^{-\tau(\hat{H}_0 - \mu\hat{N})} S(\tau)$$

$$= S^{\dagger}(\tau) e^{-\tau(\hat{H}_0 - \mu\hat{N})}$$

$$\frac{\partial S(\tau)}{\partial \tau} = -\hat{H}_{int} S(\tau) \rightarrow S(\tau) = T_{\tau} e^{-\int_0^{\tau} \hat{H}_{int}(\tau') d\tau'}$$

$$G(\tau_1, \tau_2) = - \frac{1}{\Xi} \left( e^{-\beta(\hat{H}_0 - \mu\hat{N})} S(\beta) S^{-1}(\tau_1) \underbrace{e^{\tau_1(\hat{H}_0 - \mu\hat{N})} \hat{\psi}(\vec{r}_1)}_{\hat{\psi}_0(\vec{r}_1, \tau_1)} e^{-\tau_1(\hat{H}_0 - \mu\hat{N})} S(\tau_1) S^{-1}(\tau_2) \underbrace{e^{\tau_2(\hat{H}_0 - \mu\hat{N})} \hat{\psi}^\dagger(\vec{r}_2)}_{\hat{\psi}_0^\dagger(\vec{r}_2, \tau_2)} e^{-\tau_2(\hat{H}_0 - \mu\hat{N})} S(\tau_2) \right)$$

$$= - \frac{1}{\Xi} \text{tr} \left\{ e^{-\beta(\hat{H}_0 - \mu\hat{N})} S(\beta) S^{-1}(\tau_1) \hat{\psi}_0(\vec{r}_1, \tau_1) S(\tau_1) S^{-1}(\tau_2) \hat{\psi}_0^\dagger(\vec{r}_2, \tau_2) S(\tau_2) \right\}$$

$$S(\tau_1, \tau_2) = T_\tau e^{-\int_{\tau_1}^{\tau_2} \hat{H}_{\text{int}}(\tau') d\tau'} = S(\tau_1) S^{-1}(\tau_2)$$

$$\Delta = - \frac{1}{\Xi} \text{tr} \left\{ e^{-\beta(\hat{H}_0 - \mu\hat{N})} S(\beta, \tau_1) \hat{\psi}_0(\vec{r}_1, \tau_1) S(\tau_1, \tau_2) \hat{\psi}_0^\dagger(\vec{r}_2, \tau_2) S(\tau_2, 0) \right\}$$

$$\Delta = - \frac{1}{\Xi} \text{tr} \left\{ e^{-\beta(\hat{H}_0 - \mu\hat{N})} T_\tau \left( \hat{\psi}_0(\vec{r}_1, \tau_1) \hat{\psi}_0^\dagger(\vec{r}_2, \tau_2) S(\beta, 0) \right) \right\}$$

↑ 包含了相互作用

$$\Xi = \text{tr} \left( e^{-\beta(\hat{H} - \mu\hat{N})} \right)$$

$$= \text{tr} \left( e^{-\beta(\hat{H}_0 - \mu\hat{N})} S(\beta) \right)$$

$$\langle\langle \dots \rangle\rangle = \frac{\text{tr}(\dots e^{-\beta(\hat{H}_0 - \mu\hat{N})})}{\text{tr}(e^{-\beta(\hat{H}_0 - \mu\hat{N})})}$$

$$G(\vec{r}_1, \tau_1, \vec{r}_2, \tau_2) = \frac{- \langle\langle T_\tau \hat{\psi}_0(\vec{r}_1, \tau_1) \hat{\psi}_0^\dagger(\vec{r}_2, \tau_2) S(\beta) \rangle\rangle_0}{\langle\langle S(\beta) \rangle\rangle_0}$$

起点

以电荷相互作用为例作微扰展开

$$G(\vec{p}, \tau) = - \frac{\langle\langle T_\tau S(\beta) C_{\vec{p}}(\tau) C_{\vec{p}}^\dagger(0) \rangle\rangle_0}{\langle\langle S(\beta) \rangle\rangle_0}$$

计算时

→ 贡献“泡泡图”，当作 1

相互作用写到动表象下

$$\hat{H}_{\text{int}} = \sum_{\vec{q}} M_{\vec{q}} (\hat{b}_{\vec{q}} + \hat{b}_{-\vec{q}}^\dagger) \sum_{\vec{k}} \hat{C}_{\vec{k}+\vec{q}}^\dagger \hat{C}_{\vec{k}}$$

$$S(\beta) = e^{-\int_0^\beta dt' H_{\text{int}}(t')} \approx 1 - \int_0^\beta dt' H_{\text{int}}(t') + \frac{1}{2} \int_0^\beta dt_1 \int_0^\beta dt_2 \hat{H}_{\text{int}}(t_1) \hat{H}_{\text{int}}(t_2)$$

- 一阶项没有贡献:  $\sim \langle \langle T_\tau (b+b^\dagger) c c^\dagger \rangle \rangle \Rightarrow \langle \langle b \rangle \rangle = 0$

$$\text{即 } G^{(1)} = 0$$

$$G^{(2)}(\vec{p}, \tau) = \frac{1}{2} \int_0^\beta dt_1 \int_0^\beta dt_2 \langle \langle T_\tau \hat{C}_{\vec{p}}(\tau) \hat{H}_{\text{int}}(t_1) \hat{H}_{\text{int}}(t_2) \hat{C}_{\vec{p}}^\dagger(0) \rangle \rangle$$

$$= \frac{1}{2} \int_0^\beta dt_1 \int_0^\beta dt_2 \langle \langle T_\tau \hat{C}_{\vec{p}}(\tau) \sum_{\vec{k}, \vec{q}} M_{\vec{q}} (\hat{b}_{\vec{q}}(t_1) + \hat{b}_{-\vec{q}}^\dagger(t_1)) \hat{C}_{\vec{k}+\vec{q}}^\dagger(t_1) \hat{C}_{\vec{k}}(t_1)$$

$$\times \sum_{\vec{k}', \vec{q}'} M_{\vec{q}'} (\hat{b}_{\vec{q}'}(t_2) + \hat{b}_{-\vec{q}'}^\dagger(t_2)) \hat{C}_{\vec{k}'+\vec{q}'}^\dagger(t_2) \hat{C}_{\vec{k}'}(t_2) \hat{C}_{\vec{p}}^\dagger(0) \rangle \rangle$$

$$= \frac{1}{2} \sum_{\vec{k}, \vec{q}} \sum_{\vec{k}', \vec{q}'} \int_0^\beta dt_1 \int_0^\beta dt_2 M_{\vec{q}} M_{\vec{q}'} \langle \langle T_\tau \hat{C}_{\vec{p}}(\tau) \hat{C}_{\vec{k}+\vec{q}}^\dagger(t_1) \hat{C}_{\vec{k}}(t_1) \hat{C}_{\vec{k}'+\vec{q}'}^\dagger(t_2) \hat{C}_{\vec{k}'}(t_2) \hat{C}_{\vec{p}}^\dagger(0) \rangle \rangle$$

$$\times \langle \langle (\hat{b}_{\vec{q}}(t_1) + \hat{b}_{-\vec{q}}^\dagger(t_1)) (\hat{b}_{\vec{q}'}(t_2) + \hat{b}_{-\vec{q}'}^\dagger(t_2)) \rangle \rangle$$

Wick 定理

$$\langle \langle AB^\dagger CD^\dagger \rangle \rangle = \langle \langle AB^\dagger \rangle \rangle \langle \langle CD^\dagger \rangle \rangle + \langle \langle AD^\dagger \rangle \rangle \langle \langle BC \rangle \rangle$$

$$\langle \langle T_\tau \hat{C}_{\vec{p}}(\tau) \hat{C}_{\vec{k}+\vec{q}}^\dagger(t_1) \hat{C}_{\vec{k}}(t_1) \hat{C}_{\vec{k}'+\vec{q}'}^\dagger(t_2) \hat{C}_{\vec{k}'}(t_2) \hat{C}_{\vec{p}}^\dagger(0) \rangle \rangle$$

$$= \langle \langle T_\tau \hat{C}_{\vec{p}}(\tau) \hat{C}_{\vec{k}+\vec{q}}^\dagger(t_1) \rangle \rangle_0 \langle \langle T_\tau \hat{C}_{\vec{k}}(t_1) \hat{C}_{\vec{k}'+\vec{q}'}^\dagger(t_2) \rangle \rangle_0 \langle \langle T_\tau \hat{C}_{\vec{k}'}(t_2) \hat{C}_{\vec{p}}^\dagger(0) \rangle \rangle_0$$

$$+ \langle \langle T_\tau \hat{C}_{\vec{p}}(\tau) \hat{C}_{\vec{k}'+\vec{q}'}^\dagger(t_2) \rangle \rangle_0 \langle \langle T_\tau \hat{C}_{\vec{k}}(t_1) \hat{C}_{\vec{p}}^\dagger(0) \rangle \rangle_0 \langle \langle T_\tau \hat{C}_{\vec{k}+\vec{q}}^\dagger(t_2) \hat{C}_{\vec{k}'}(t_2) \rangle \rangle_0$$

$$+ \langle \langle T_\tau \hat{C}_{\vec{p}}(\tau) \hat{C}_{\vec{p}}^\dagger(0) \rangle \rangle_0 \langle \langle T_\tau \hat{C}_{\vec{k}+\vec{q}}^\dagger(t_1) \hat{C}_{\vec{k}'}(t_2) \rangle \rangle_0 \langle \langle T_\tau \hat{C}_{\vec{k}}(t_1) \hat{C}_{\vec{k}'+\vec{q}'}^\dagger(t_2) \rangle \rangle_0$$

$$+ \langle \langle T_\tau \hat{C}_{\vec{p}}(\tau) \hat{C}_{\vec{k}'+\vec{q}'}^\dagger(t_2) \rangle \rangle_0 \langle \langle T_\tau \hat{C}_{\vec{k}+\vec{q}}^\dagger(t_1) \hat{C}_{\vec{k}'}(t_2) \rangle \rangle_0 \langle \langle T_\tau \hat{C}_{\vec{k}}(t_1) \hat{C}_{\vec{p}}^\dagger(0) \rangle \rangle_0$$

$$+ \langle \langle T_\tau \hat{C}_{\vec{p}}(\tau) \hat{C}_{\vec{k}'+\vec{q}'}^\dagger(t_2) \rangle \rangle_0 \langle \langle T_\tau \hat{C}_{\vec{k}+\vec{q}}^\dagger(t_1) \hat{C}_{\vec{k}'}(t_2) \rangle \rangle_0 \langle \langle T_\tau \hat{C}_{\vec{p}}^\dagger(0) \hat{C}_{\vec{k}}(t_1) \rangle \rangle_0$$

$$+ \langle \langle T_\tau \hat{C}_{\vec{p}}(\tau) \hat{C}_{\vec{p}}^\dagger(0) \rangle \rangle_0 \langle \langle T_\tau \hat{C}_{\vec{k}+\vec{q}}^\dagger(t_1) \hat{C}_{\vec{k}}(t_1) \rangle \rangle_0 \langle \langle T_\tau \hat{C}_{\vec{k}'+\vec{q}'}^\dagger(t_2) \hat{C}_{\vec{k}'}(t_2) \rangle \rangle_0$$

$$= -G_{\vec{P}}^0(\tau, \tau_1) G_{\vec{R}}^0(\tau_1, \tau_2) G_{\vec{R}}^0(\tau_2, 0) \delta_{\vec{P}, \vec{R}+\vec{q}} \delta_{\vec{R}, \vec{R}+\vec{q}} \delta_{\vec{R}, \vec{P}} \quad (1)$$

$$+ G_{\vec{P}}^0(\tau, \tau_1) G_{\vec{R}}^0(\tau_1, 0) G_{\vec{R}}^0(\tau_2, \tau_2) \delta_{\vec{P}, \vec{R}+\vec{q}} \delta_{\vec{R}, \vec{P}} \delta_{\vec{R}+\vec{q}, \vec{R}} \quad (2)$$

$$+ G_{\vec{P}}^0(\tau, 0) G_{\vec{R}}^0(\tau_2, \tau_1) G_{\vec{R}}^0(\tau_1, \tau_2) \delta_{\vec{R}, \vec{R}+\vec{q}} \delta_{\vec{R}, \vec{R}+\vec{q}} \quad (3)$$

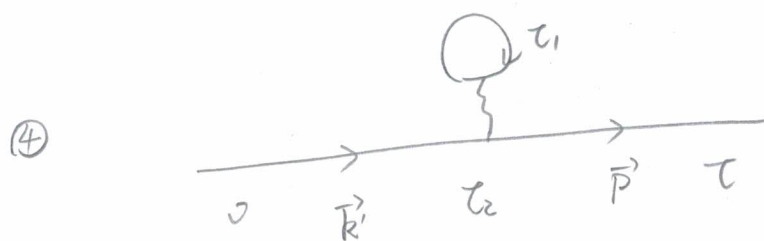
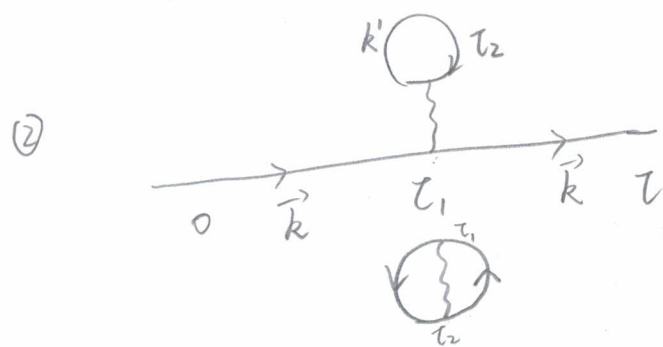
$$+ G_{\vec{P}}^0(\tau, \tau_2) G_{\vec{R}}^0(\tau_1, \tau_1) G_{\vec{R}}^0(\tau_2, 0) \delta_{\vec{P}, \vec{R}+\vec{q}} \delta_{\vec{R}, \vec{R}+\vec{q}} \delta_{\vec{R}, \vec{P}} \quad (4)$$

$$- G_{\vec{P}}^0(\tau, \tau_2) G_{\vec{R}}^0(\tau_2, \tau_1) G_{\vec{R}}^0(\tau_1, 0) \delta_{\vec{P}, \vec{R}+\vec{q}} \delta_{\vec{R}, \vec{R}+\vec{q}} \delta_{\vec{R}, \vec{P}} \quad (5)$$

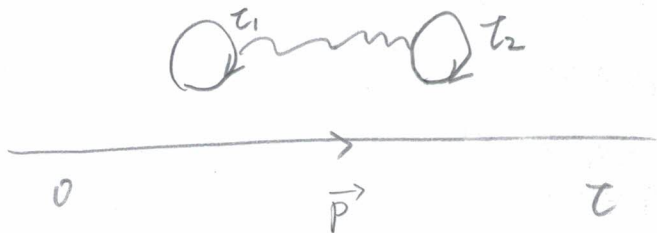
$$- G_{\vec{P}}^0(\tau, 0) G_{\vec{R}}^0(\tau_1, \tau_1) G_{\vec{R}}^0(\tau_2, \tau_2) \delta_{\vec{R}, \vec{R}+\vec{q}} \delta_{\vec{R}, \vec{R}+\vec{q}} \quad (6)$$

把  $\langle\langle (\hat{b}_{\vec{q}}(\tau_1) + \hat{b}_{-\vec{q}}^+(\tau_1)) (\hat{b}_{\vec{q}}(\tau_2) + \hat{b}_{-\vec{q}}^+(\tau_2)) \rangle\rangle$  记为  $V(\tau_1, \tau_2)$

则  $\langle\langle \hat{C}_{\vec{P}}(\tau) \hat{C}_{\vec{R}+\vec{q}}(\tau_1) \hat{C}_{\vec{R}}(\tau_1) \hat{C}_{\vec{R}+\vec{q}}^+(\tau_2) \hat{C}_{\vec{R}}(\tau_2) \hat{C}_{\vec{P}}^+(0) \rangle\rangle$  的 6 个费曼图



(6)



只计入第①和第⑤的贡献

$$\begin{aligned}
 G^{(2)}(\vec{p}, \tau) = & \left( \frac{1}{2} \int_{\vec{q}} \int_0^\beta dt_1 \int_0^\beta dt_2 M_q M_{-q} G_{\vec{p}}^0(\tau, t_1) G_{\vec{p}-\vec{q}}^0(t_1, t_2) G_{\vec{p}}^0(t_2, 0) \right. \\
 & \left. - \frac{1}{2} \int_{\vec{q}} \int_0^\beta dt_1 \int_0^\beta dt_2 M_q M_{-q} G_{\vec{p}}^0(\tau, t_2) G_{\vec{p}+\vec{q}}^0(t_2, t_1) G_{\vec{p}}^0(t_1, 0) \right) \\
 & \times \langle\langle (\hat{b}_{\vec{q}}^-(t_1) + \hat{b}_{-\vec{q}}^+(t_1)) (\hat{b}_{-\vec{q}}^-(t_2) + \hat{b}_{\vec{q}}^+(t_2)) \rangle\rangle.
 \end{aligned}$$